

# Lecture 2 - Superradiance FEL

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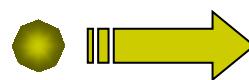
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Department of Physics

National Tsinghua University, Hsinchu, Taiwan



# 1. superradiance

# Radiation Spectral Energy of a Single Electron

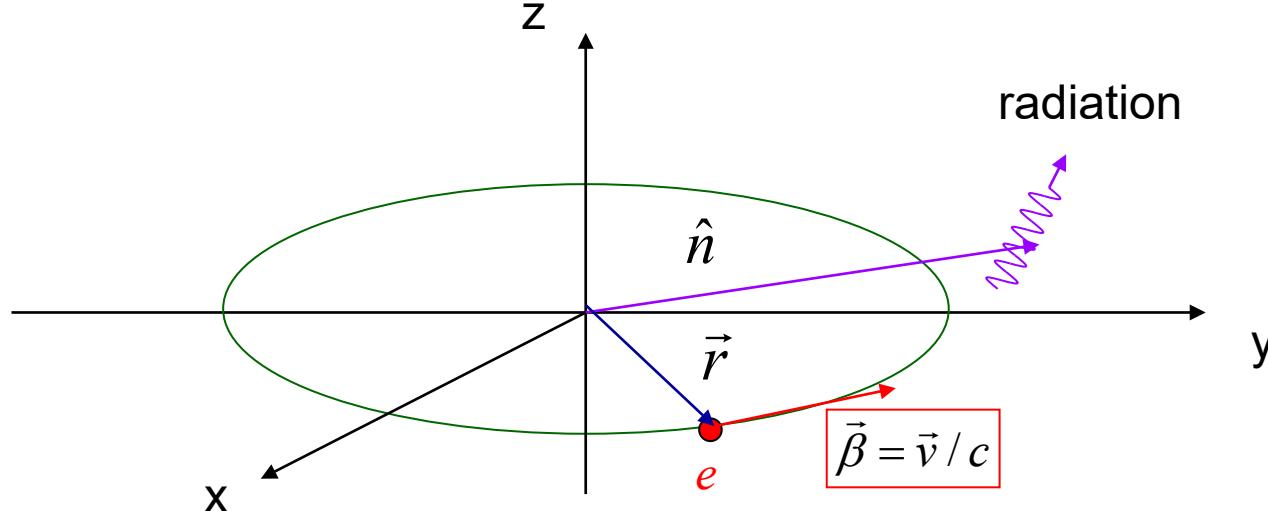


$W$ : radiation energy  
 $\omega$ : radiation frequency  
 $\mu_0$ : vacuum permeability  
 $e$ : electron charge  
 $t$ : time variable  
 $\Omega$ : solid angle

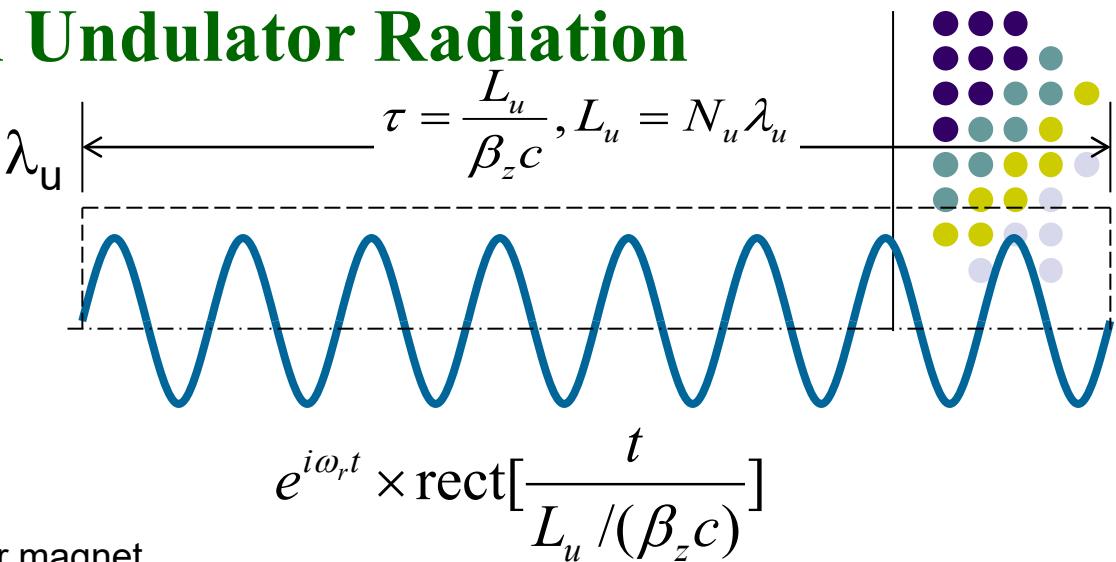
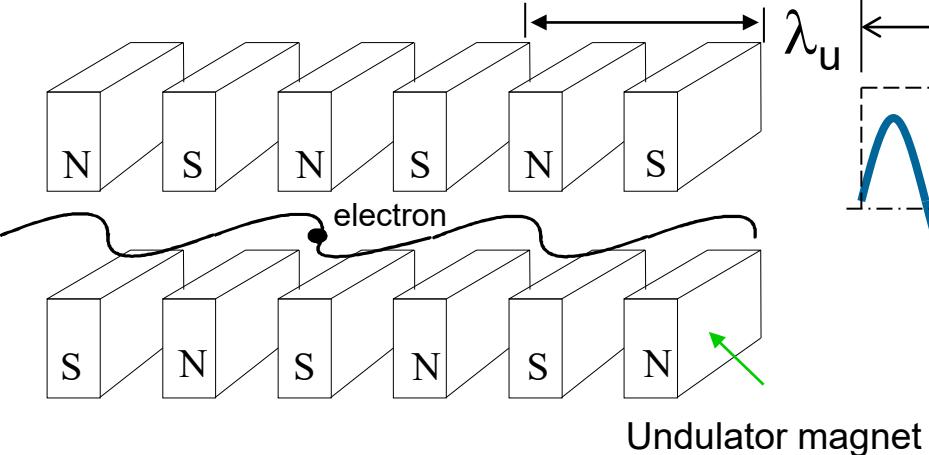
Radiation spectral energy per solid angle  $\left( \frac{d^2W}{d\omega d\Omega} \right)_1$

Single-electron Radiation spectral energy

$$\left( \frac{dW}{d\omega} \right)_1 = \frac{\mu_0 c \omega^2 e^2}{16\pi^3} \int \left| \int_{-\infty}^{\infty} (e^{j\omega(t - \hat{n} \cdot \vec{r}/c)} \hat{n} \times \vec{\beta}) dt \right|^2 d\Omega$$



# Singe-electron Undulator Radiation



## Lemma

Rectangular function:  $\text{rect}[t] = \begin{cases} 1 & \text{for } |t| \leq 1/2 \\ 0 & \text{otherwise} \end{cases}$

$N_w$ : number of wiggler periods  
 $L_w$ : wiggler length  
 $\tau$ : radiation time  
 $\omega_r$ : resonant radiation frequency

Fourier Transform  $\{\text{rect}[t]\} = \frac{\sin(\omega/2)}{\omega/2} = \text{sinc}(f)$

So,

Fourier Transform  $\{e^{i\omega_r t} \times \text{rect}\left[\frac{t}{L_u / (\beta_z c)}\right]\}$

$$\propto \frac{\sin[N_w(\omega/\omega_r - 1)]}{N_w(\omega/\omega_r - 1)}$$

Synchronous radiation frequency

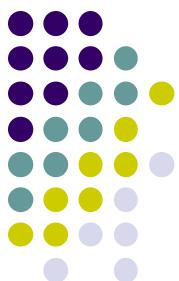
$$\omega_r = 2\pi c \left[ \frac{1 + a_u^2}{2\gamma^2} \lambda_u \right]^{-1}$$

Undulator parameter

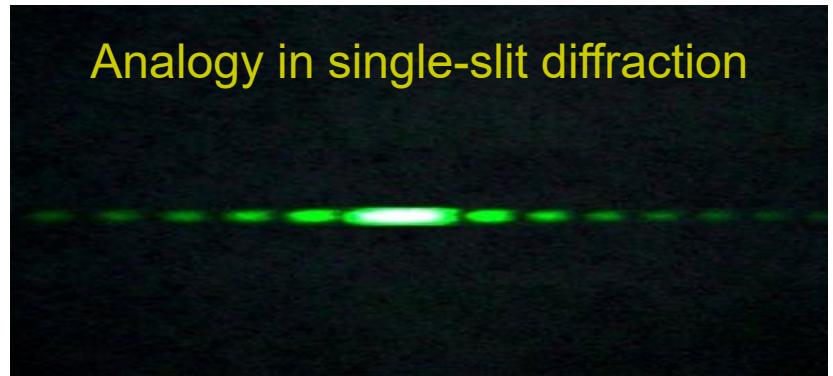
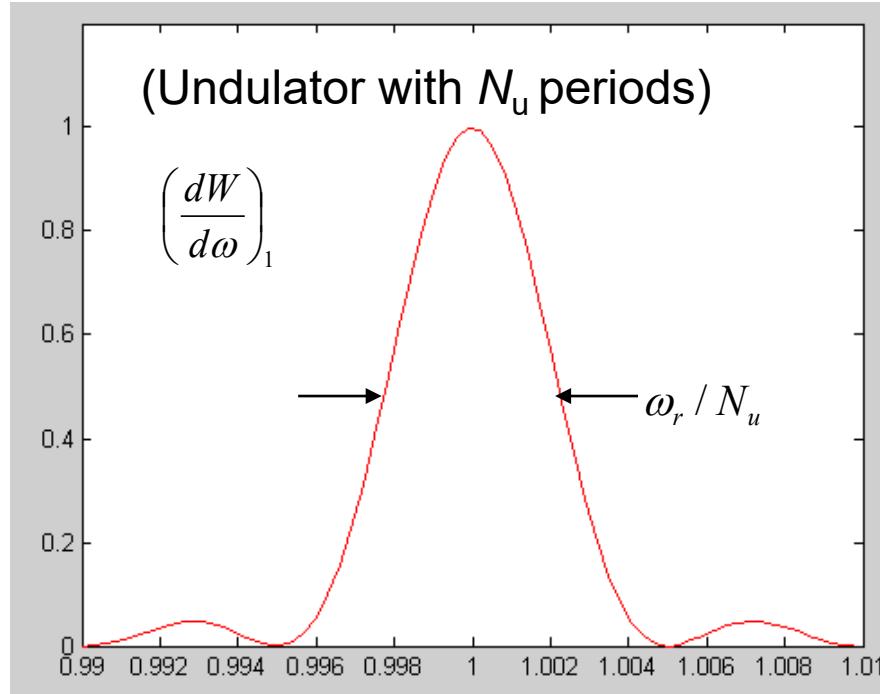
4

$$a_u = 0.093 B_{rms} (\text{kgauss}) \times \lambda_u (\text{cm})$$

**Spectral energy**  $\left(\frac{dW}{d\omega}\right)_1 \propto \{\text{sinc}[N_w(\omega/\omega_r - 1)/\pi]\}^2$



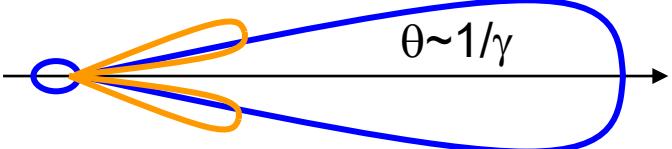
## Spectral-energy Lineshape Function



$$\omega/\omega_r$$

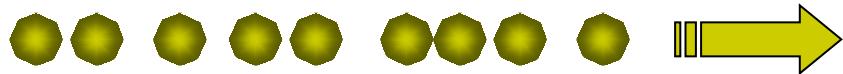
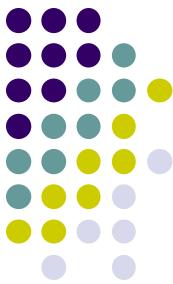
For a helical undulator

$$\left(\frac{d^2W}{d\omega d\Omega}\right)_1 = \frac{e^2 N_u^2 \gamma^2}{2\pi} \eta \frac{a_u^2}{(1 + a_u^2 + \gamma^2 \theta^2)^2} \left\{ \frac{\sin[N_u \pi (\omega/\omega_r - 1)]}{N_u \pi (\omega/\omega_r - 1)} \right\}^2$$



Solid angle  $\Delta\Omega \sim \pi / \gamma^2$      $\eta$ : Wave impedance

# Radiation from many electrons



$$\left( \frac{dW}{d\omega} \right)_N = \left( \frac{dW}{d\omega} \right)_1 \left| \sum_{i=1}^N e^{-j\phi_i} \right|^2$$

**Assume no energy exchange between electrons**

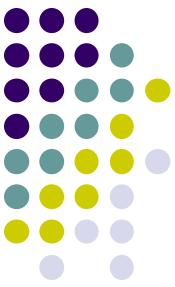
$\left( \frac{dW}{d\omega} \right)_1$  Spectral energy radiated by a single particle

$\phi_i(t, r)$  Radiation phase of  $i^{\text{th}}$  particle

$N$ : number of electrons

$$\left| \sum_{i=1}^N e^{-j\phi_i} \right|^2 = \begin{cases} \sim N & \text{for random phase } \phi_i \\ N^2 & \text{for a constant phase } \phi_i = \phi_0 \end{cases}$$

**superradiance**



## Bunching Factor

$$|M| = \left| \sum_{i=1}^N e^{-j\phi_i} \right| / N \quad 0 \leq |M| \leq 1$$

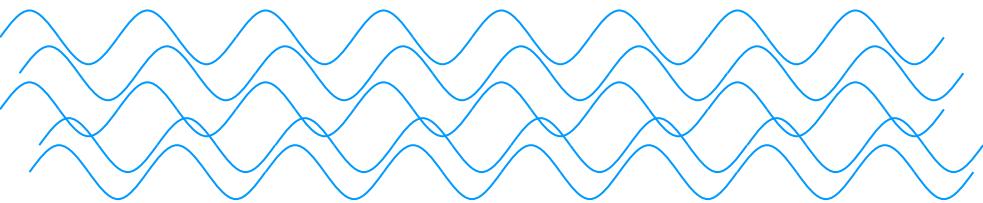
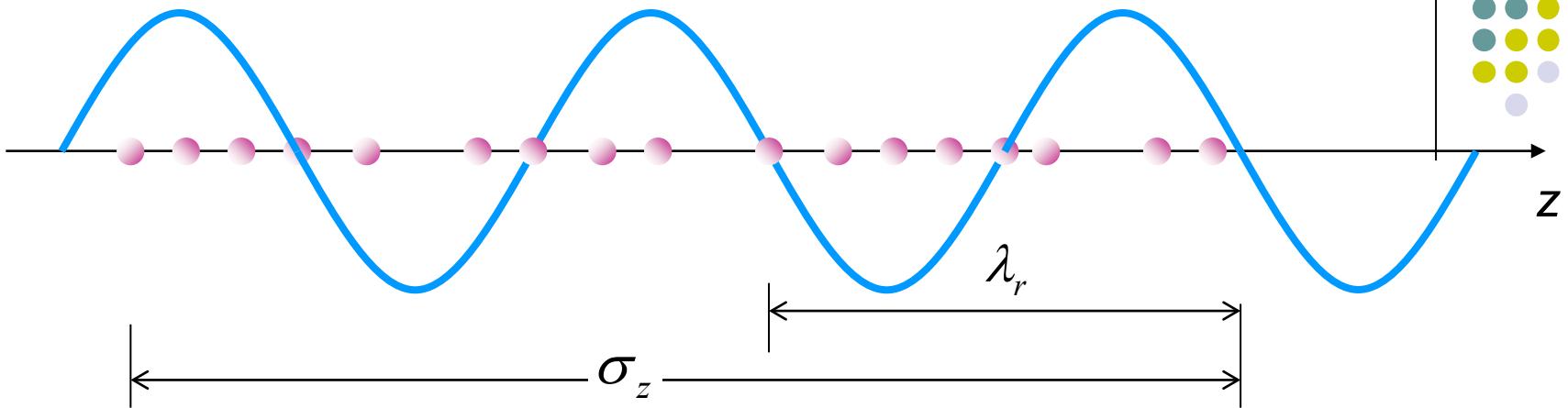
In the limit of continuous distribution with a temporal distribution function  $f(t)$ , where  $\int_{-\infty}^{\infty} f(t)dt = 1$ , the bunching factor becomes

$$|M(\omega)| \rightarrow \int_{-\infty}^{\infty} f(t)e^{-j\omega t} dt \equiv \text{ Fourier Transform of } f(t)$$

$|M(\omega_n)| = 1$  means perfect bunching at frequency  $\omega_n$ , yielding the maximum enhancement factor of  $N$  for radiation

## Incoherent Radiation

$\sigma_z \gg \lambda_r, r_i$  is random



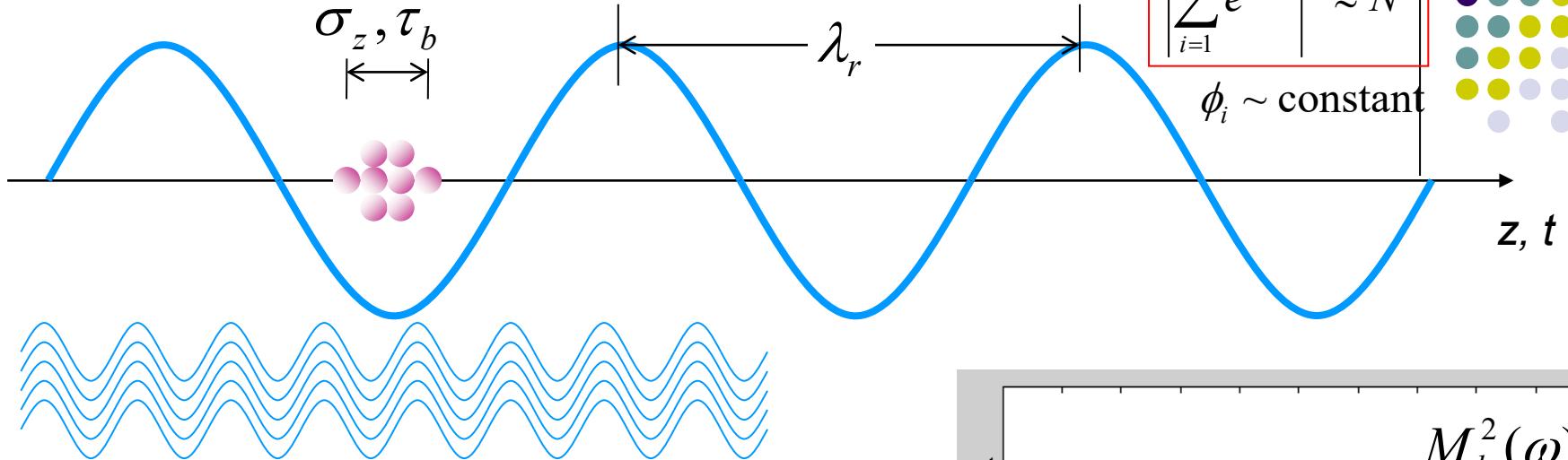
$$\left| \sum_{i=1}^N e^{-j\phi_i} \right|^2 \approx N$$

Total Spectral Energy  $\left( \frac{dW}{d\omega} \right)_{inc, N} = \boxed{N} \left( \frac{dW}{d\omega} \right)_1 \propto N$   $N$ : number of electrons

$$\left( \frac{dW}{d\omega} \right)_1 \text{ Radiation spectral energy of a single electron}$$



# Coherent Spontaneous/Synchrotron Radiation

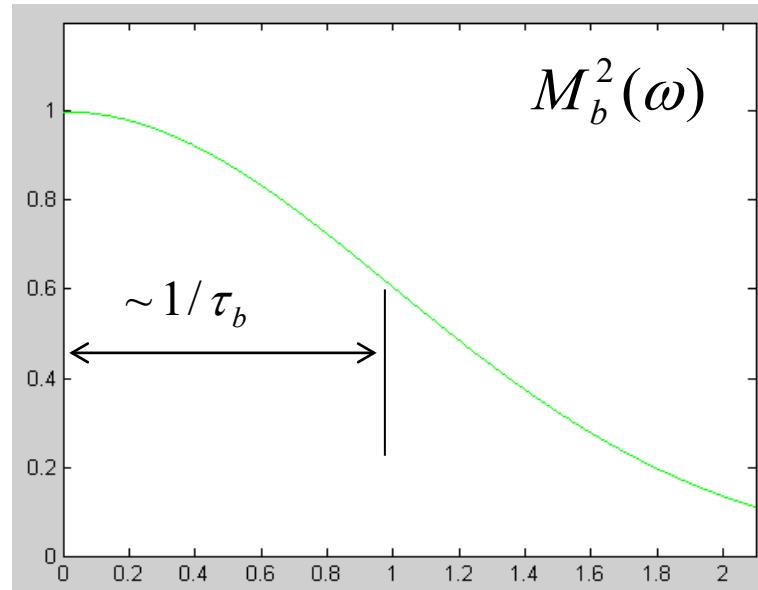


$$\text{Spectral Energy} \left( \frac{dW}{d\omega} \right)_{SR, N_b} = N_b^2 \left( \frac{dW}{d\omega} \right)_1 M_b^2(\omega)$$

$M_b(\omega)$  :Fourier transform of the bunch shape

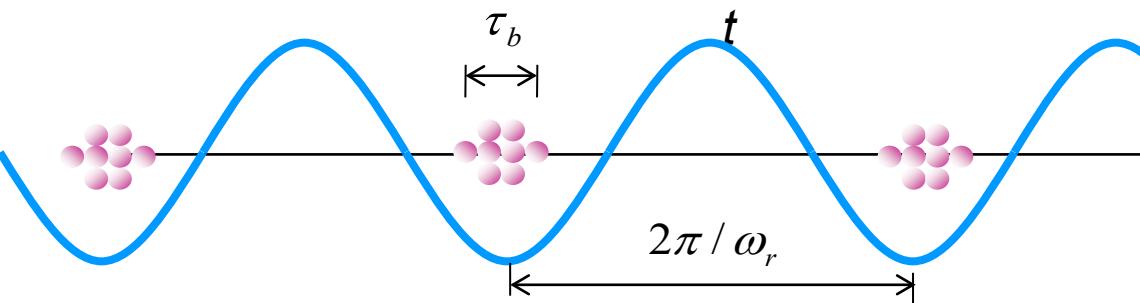
$N_b$ : number of electrons in the bunch

\* For 1 nC in 10 ps,  $N_b = 6.25 \times 10^9$ !



$$\text{For Gaussian bunch } f(t) = \frac{\exp(-t^2 / \tau_b^2)}{\sqrt{\pi\tau_b}} \Rightarrow M_b(\omega) = \exp\left(-\frac{\omega^2\tau_b^2}{4}\right)$$

# Superradiance from Periodic Bunches (comb-pulse superradiance)



Spectral Energy

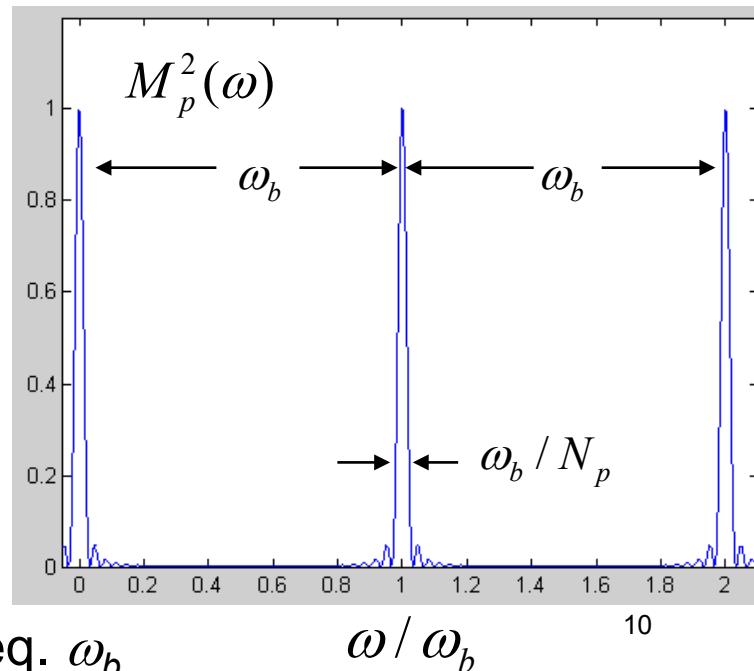
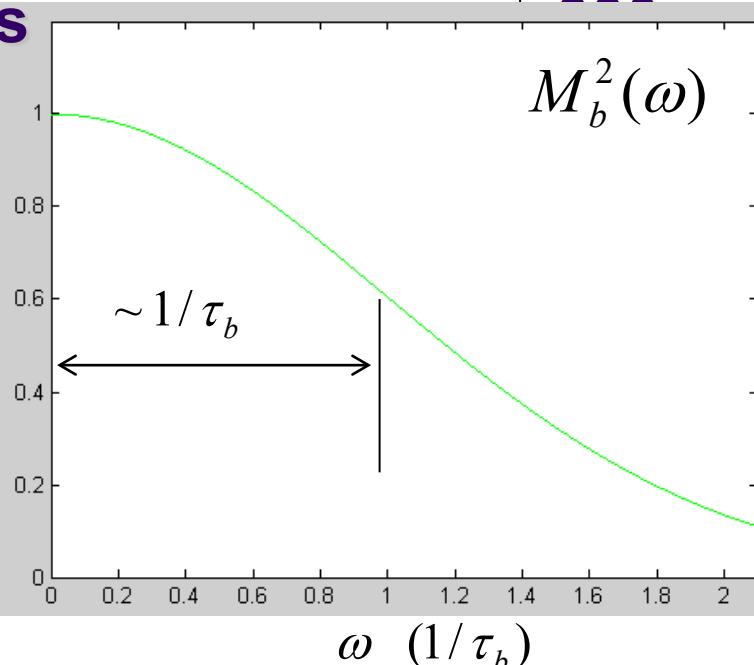
$$\left( \frac{dW}{d\omega} \right)_{SR} = \left( \frac{dW}{d\omega} \right)_1 \boxed{N_b^2 M_b^2(\omega)} \times \boxed{N_p^2 M_p^2(\omega)}$$

$M_b(\omega)$  :Fourier transform of the bunch shape

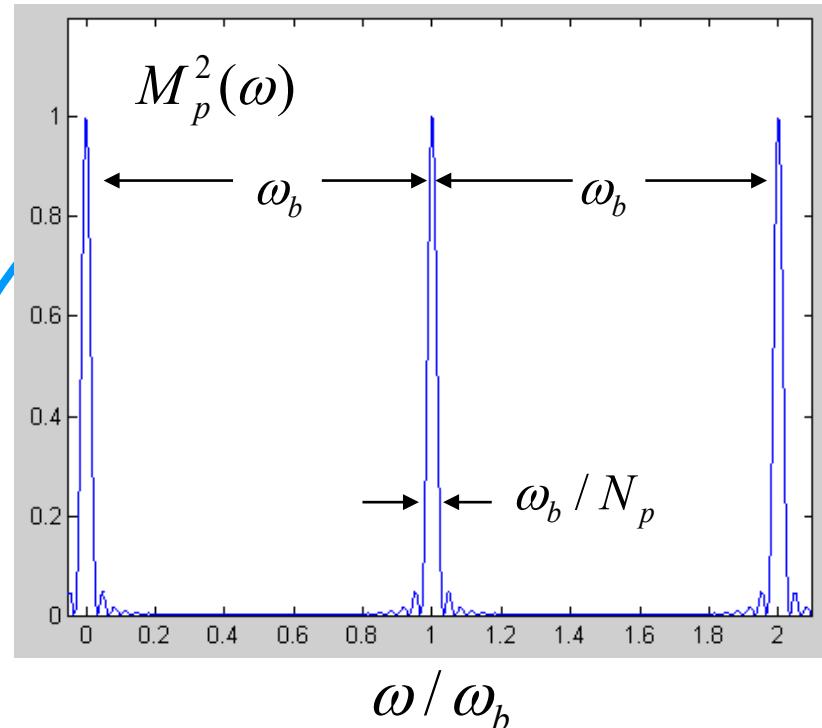
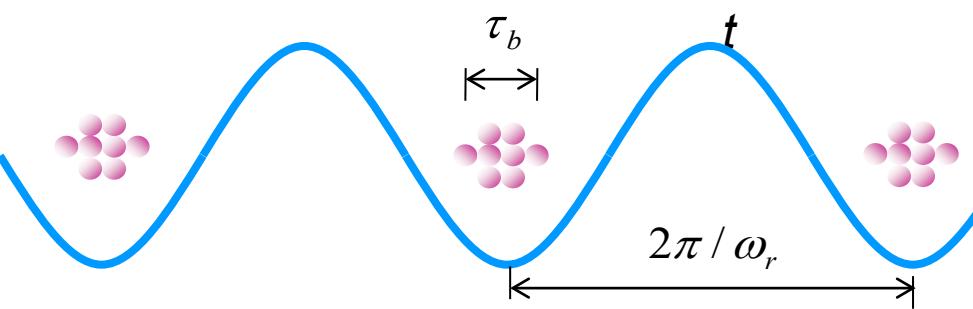
$N_b$ : number of bunched electrons

$$M_p(\omega) = \frac{\sin(N_p \pi \omega / \omega_b)}{N_p \sin(\pi \omega / \omega_b)}$$

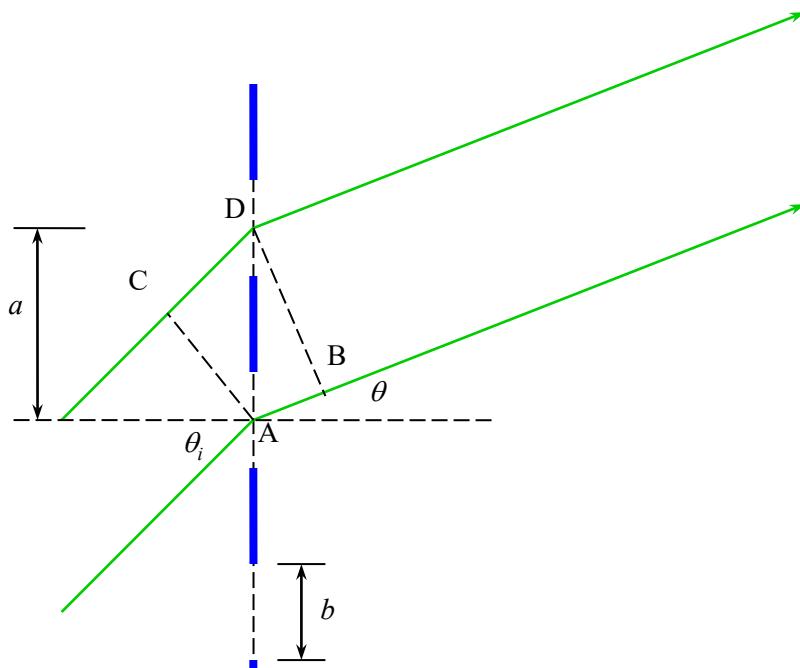
Coherent sum of  $N_p$  bunches with bunching freq.  $\omega_b$



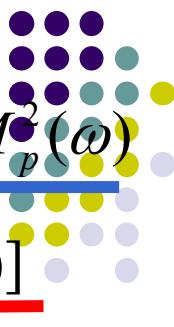
## Multi-bunch radiation



Analogy in grating diffraction



# Spectral narrowing from comb-pulse superradiance

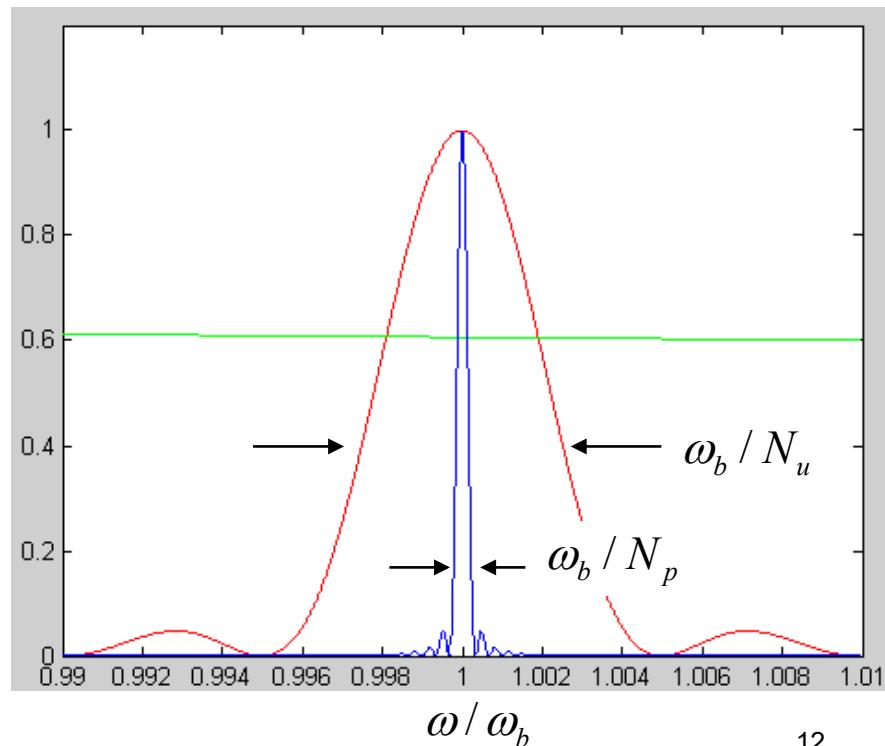
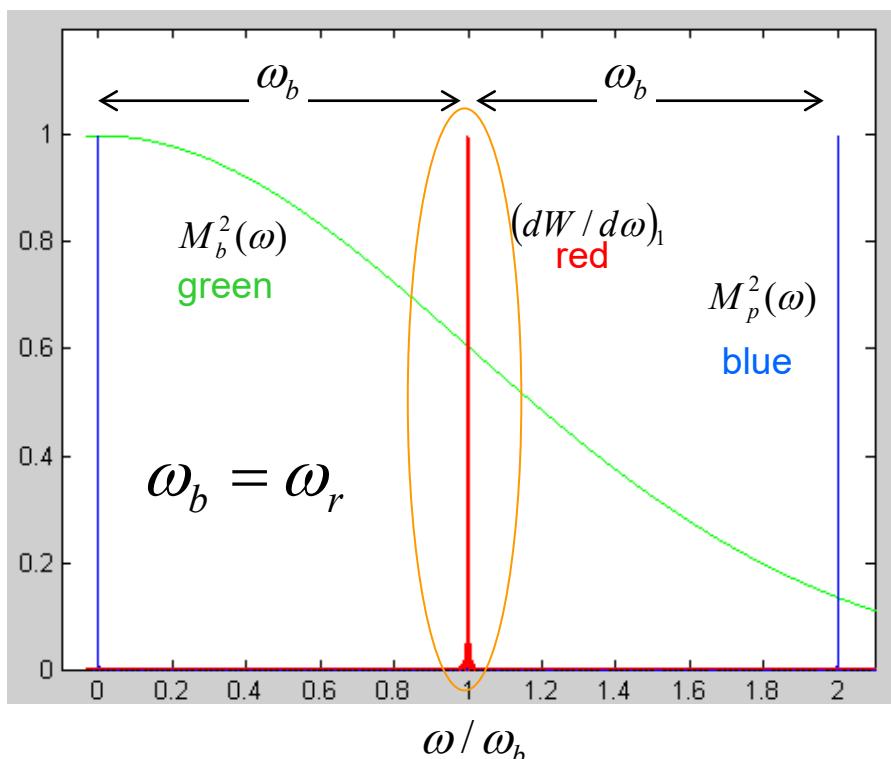


Spectral Energy  $(dW / d\omega)_{SR, N_p \times N_b} = (N_p N_b)^2 (dW / d\omega)_l M_b^2(\omega) M_p^2(\omega)$

Assume undulator radiation  $\Rightarrow (dW / d\omega)_l \propto \text{sinc}^2[N_u(\omega / \omega_r - 1)]$

$M_p(\omega) = \sin(N_p \pi \omega / \omega_b) / \sin(\pi \omega / \omega_b) / N_p$  for bunching freq. =  $\omega_b$

$|M_b(\omega)M_p(\omega)|$  is the overall *bunching factor*

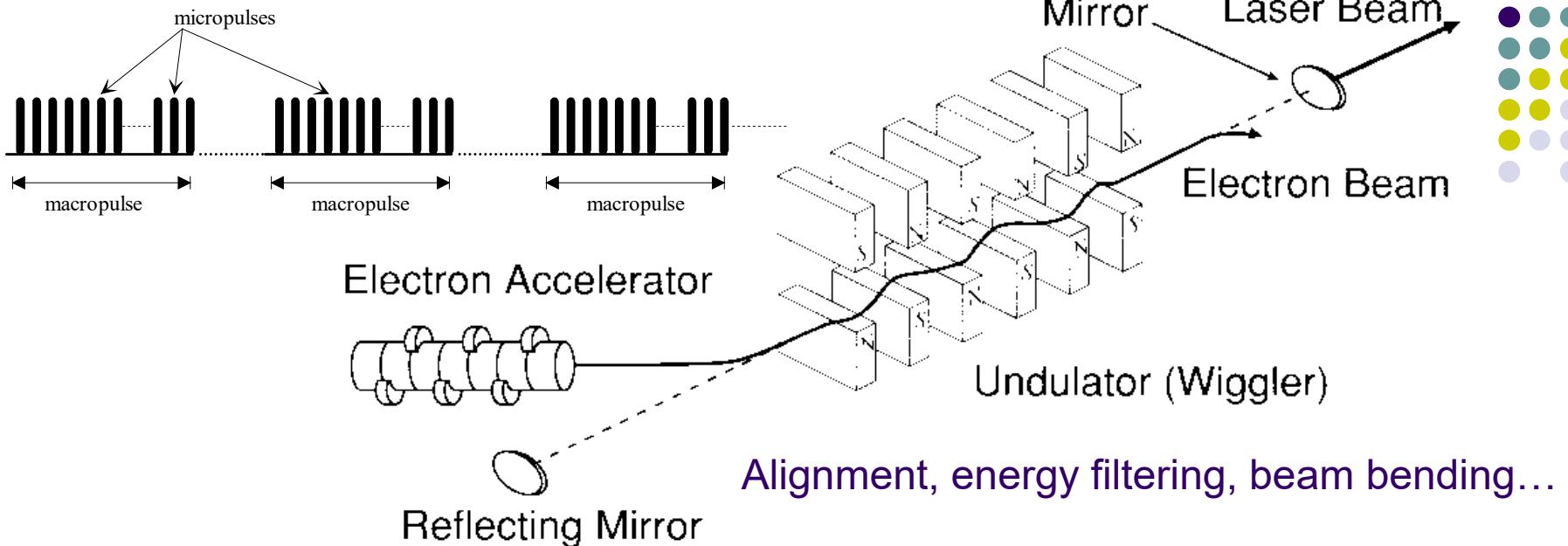


Tight bunch, large bunching factor, periodic bunching are desirable



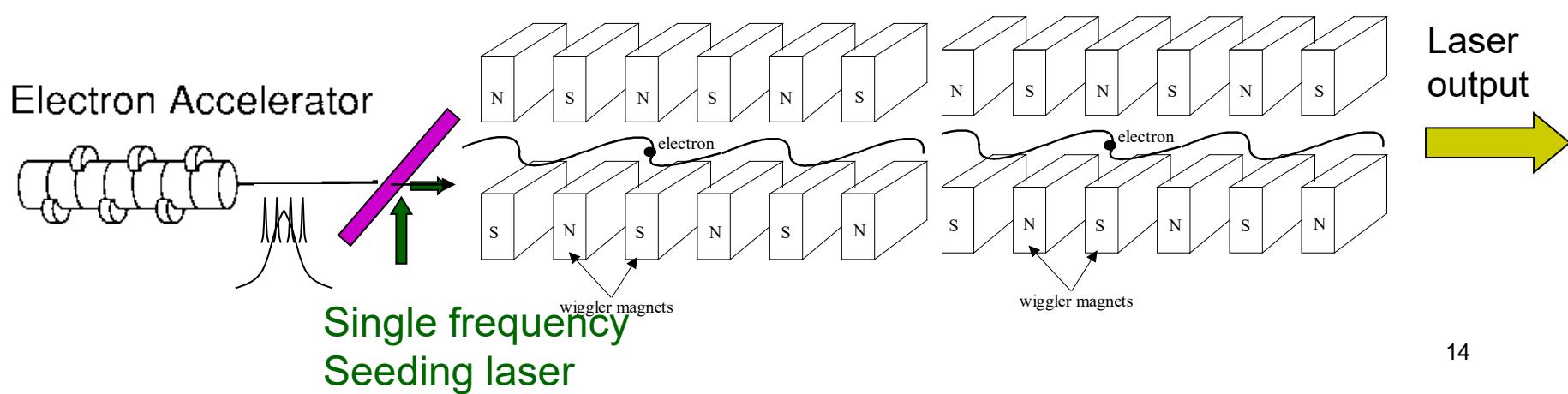
## 2. Ordinary FEL vs. Superradiance FEL

# FEL Oscillator



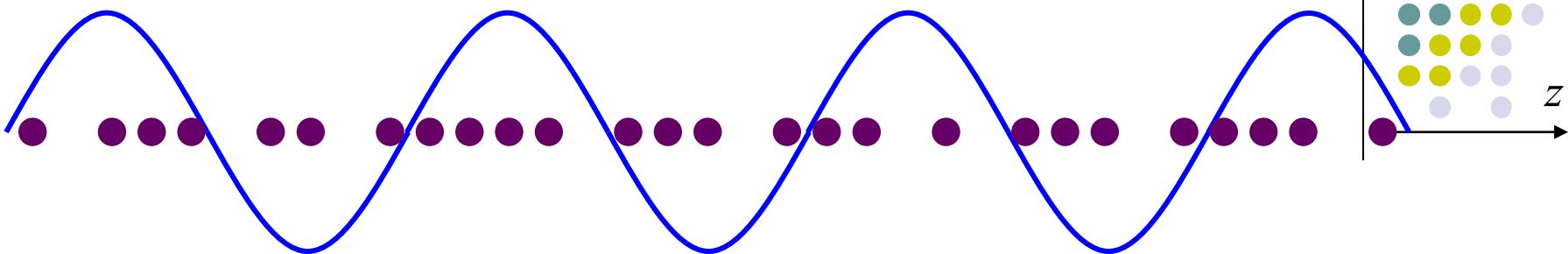
## Single-pass FEL

Long undulator for a SASE FEL



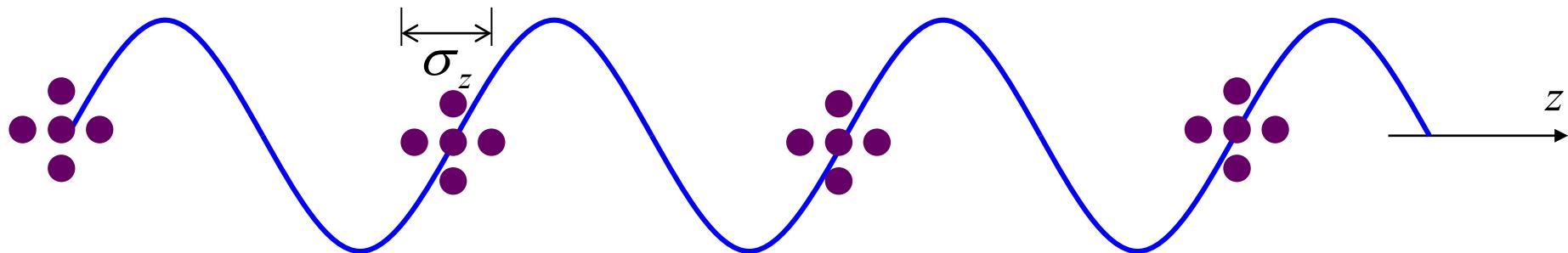
# FEL Bunching

upstream undulator



As an electron beam propagates down an FEL, electrons gain energy modulation  $\Rightarrow$  spatial modulation and are bunched to emit coherent radiation.

Down-stream undulator  $\sim$  superradiance FEL



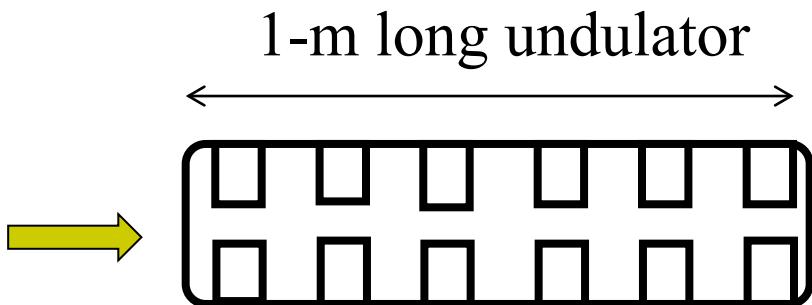
Radiation Spectral Energy

$$\left( \frac{dW}{d\omega} \right)_N = \left( \frac{dW}{d\omega} \right)_1 \left| \sum_{i=1}^N e^{-j\phi_i} \right|^2$$

Coherent Radiation

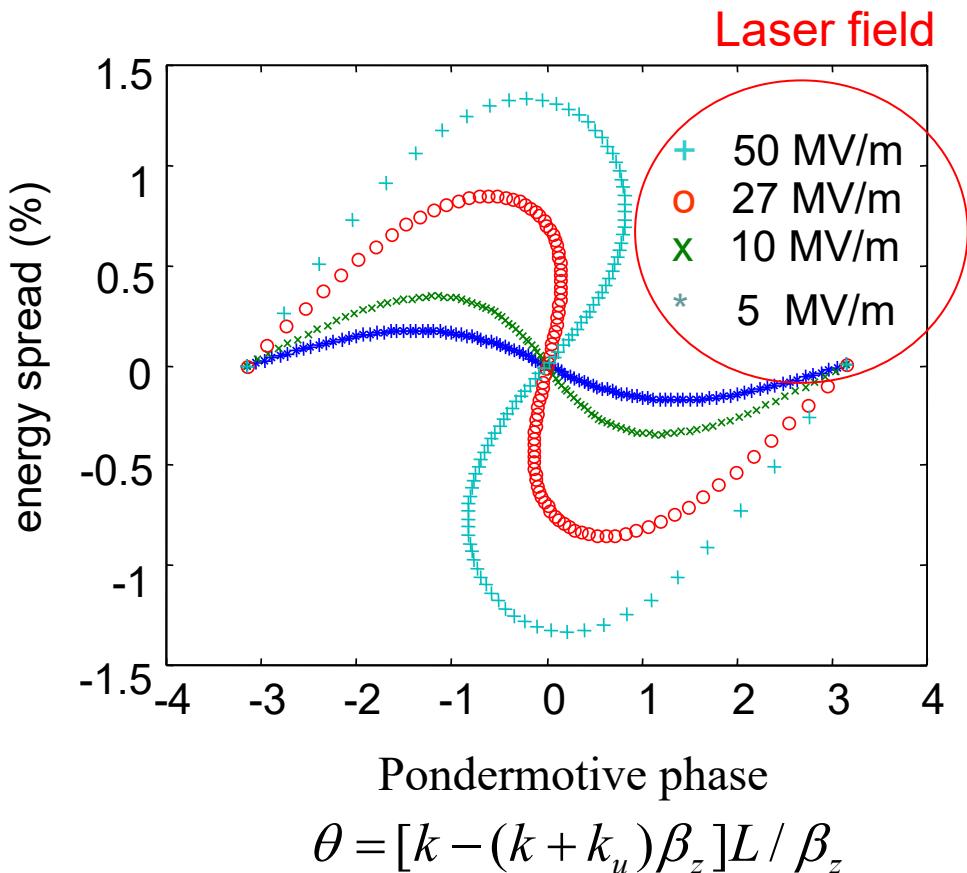
$$\sigma_z \ll \lambda, \text{ Power} \propto N^2 \quad N \sim 10^{8-10} !$$

# Energy $\Rightarrow$ Spatial Modulation in a Undulator

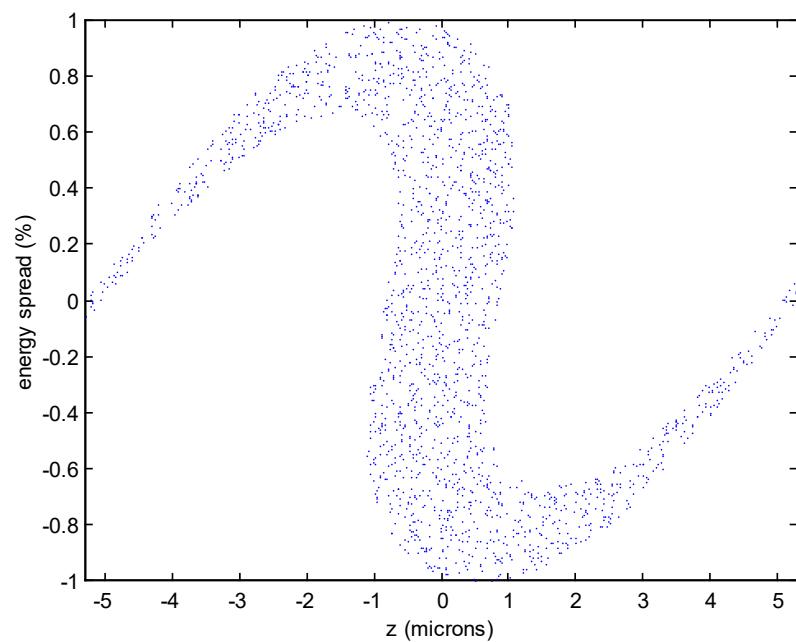



$$B_u = 2.9 \text{ kG}, \quad \lambda_u = 3.56 \text{ cm}$$

Laser @  $\lambda = 10 \mu\text{m}$



Phase matching wavelength =  $10 \mu\text{m}$



After the end of the 1-m wiggler<sub>16</sub>  
for 27MV/m laser field

# SASE FEL



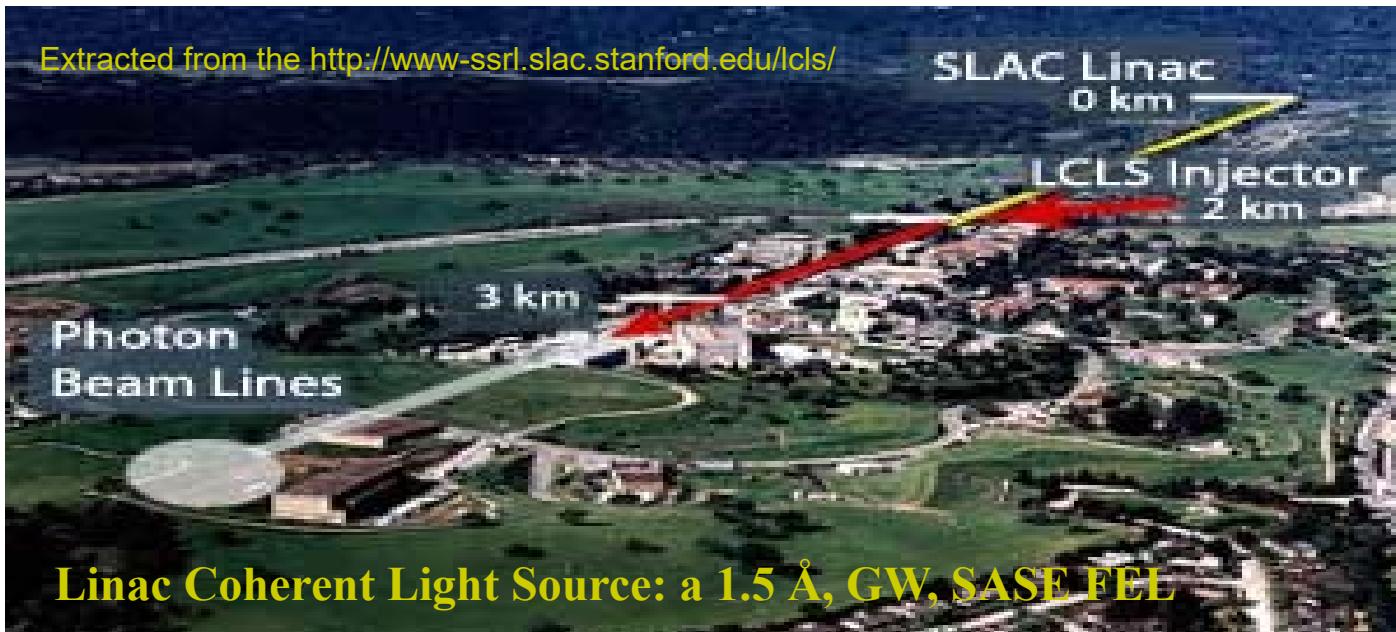
## Self-amplified Spontaneous Emission (SASE)

As electrons propagate down a long wiggler, electrons are bunched, emit coherent radiation, and amplify the radiation .

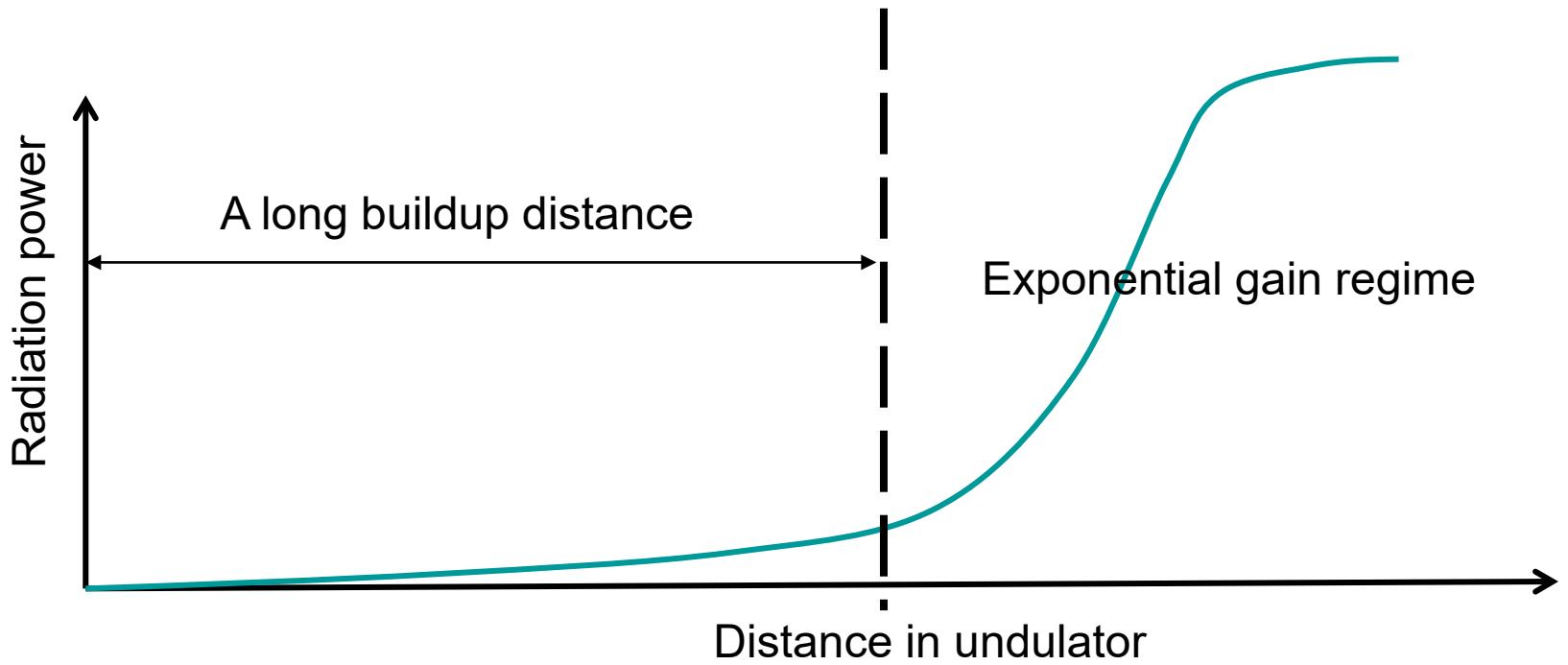
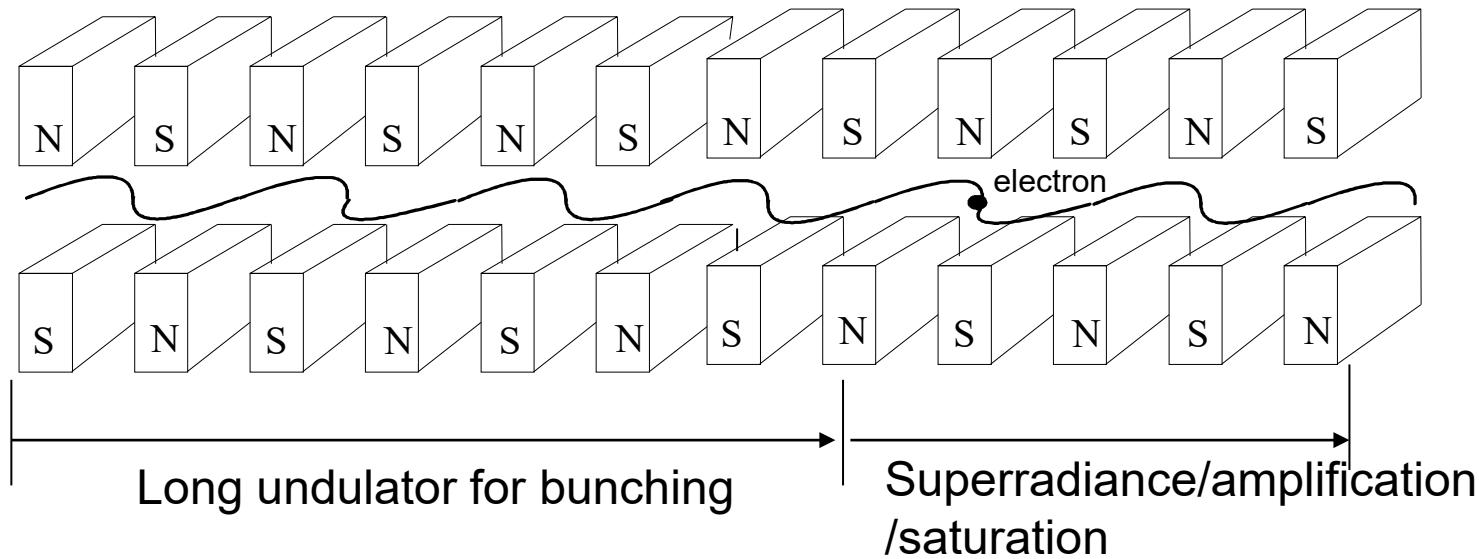
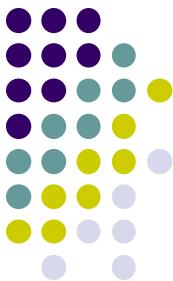
### 1-D Model

In the high-gain regime, the SASE FEL gain  $G = \frac{P(z)}{P_{in}} = \frac{1}{9} \exp(2z/L_g)$

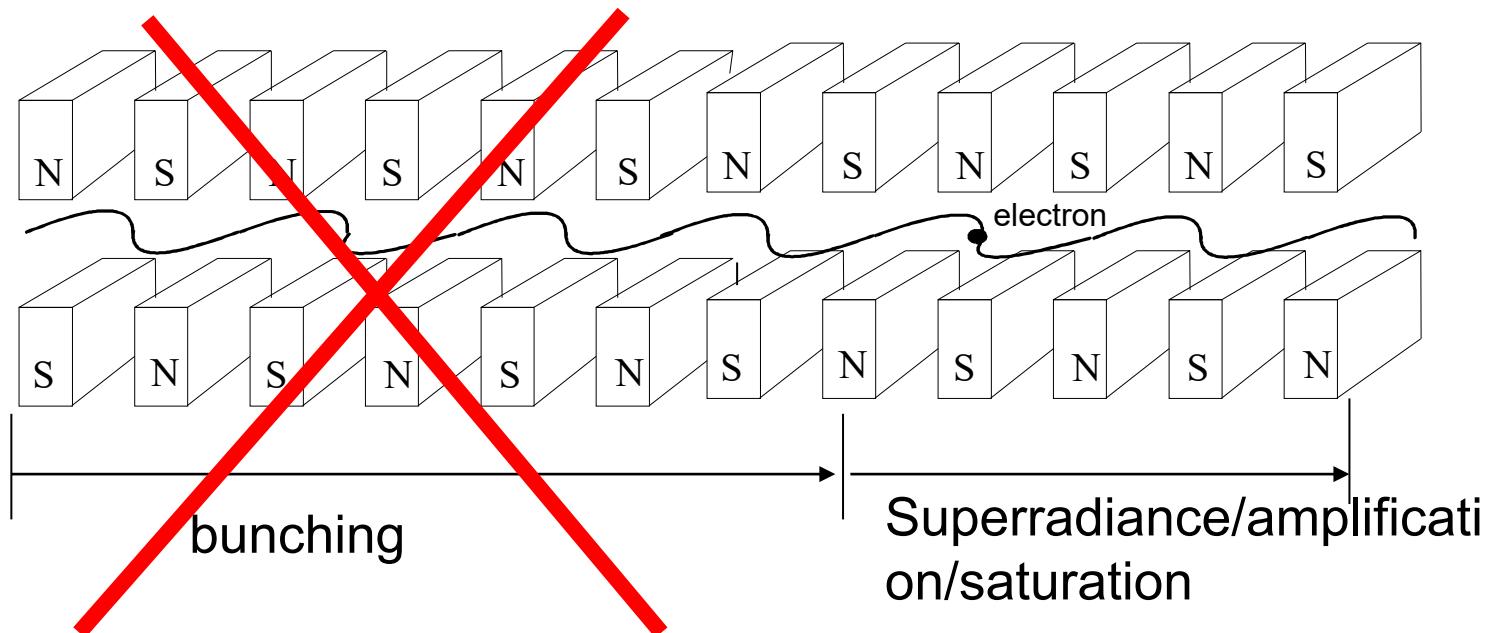
$L_g = \lambda_u / 2\sqrt{3}\pi\rho$  : gain length,  $\rho \approx \frac{0.88}{\gamma} \frac{B_u \lambda_u^{4/3} I^{1/3}}{\epsilon_n^{1/3}}$  : fundamental FEL parameter



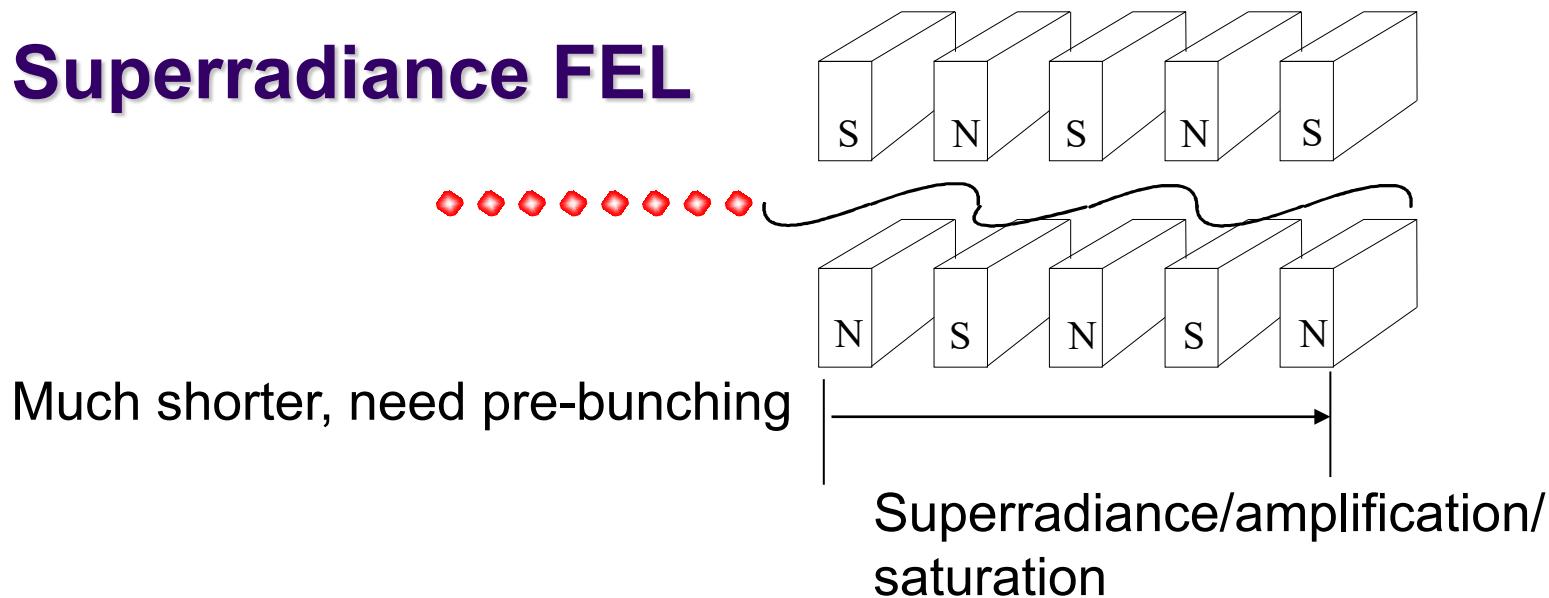
# SAFE FEL



# Conventional SAFE FEL

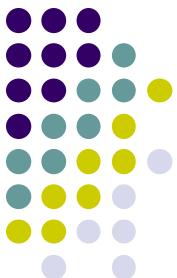


# Superradiance FEL

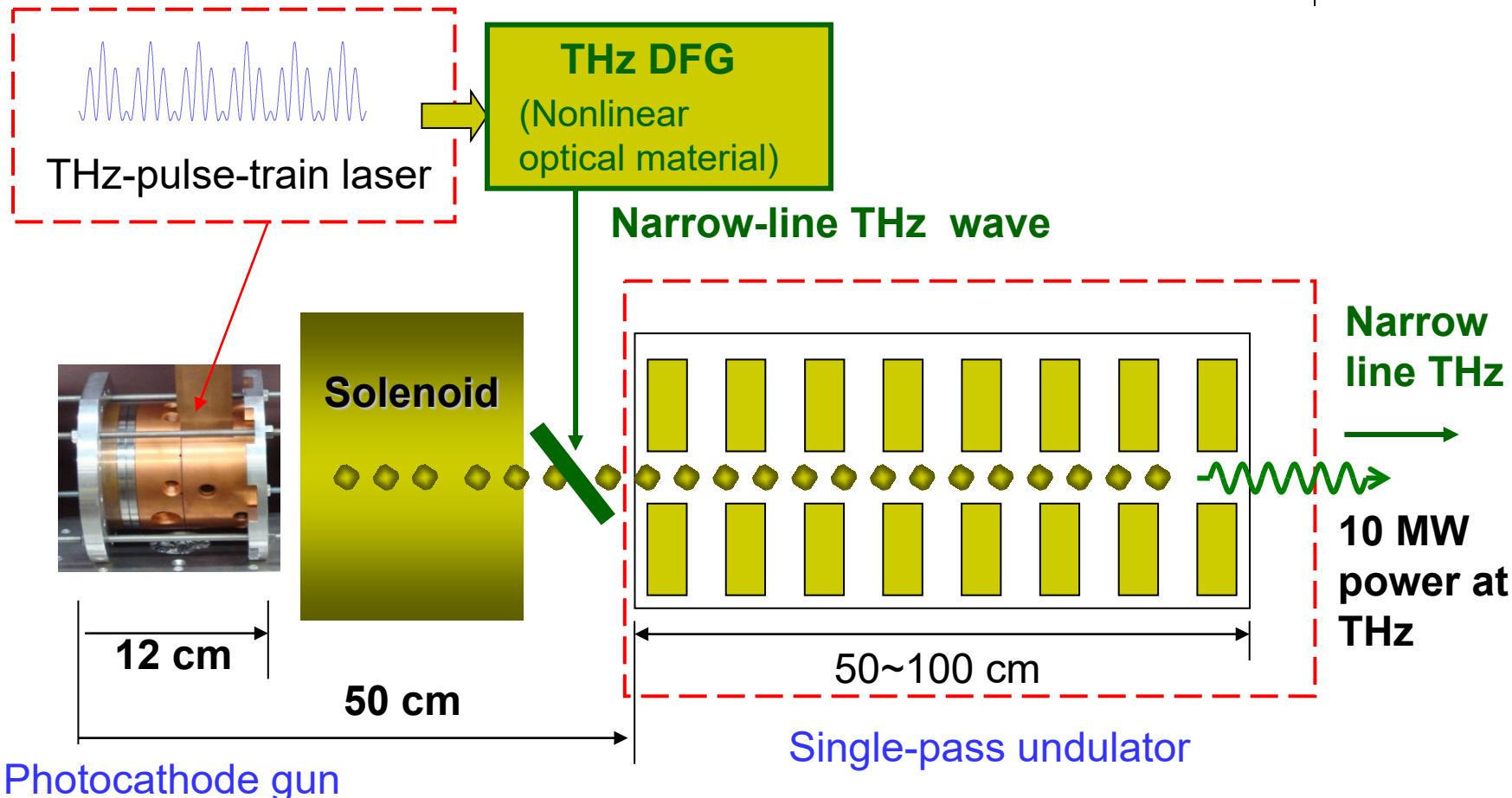




### 3. A Design Example of a Superradiance FEL



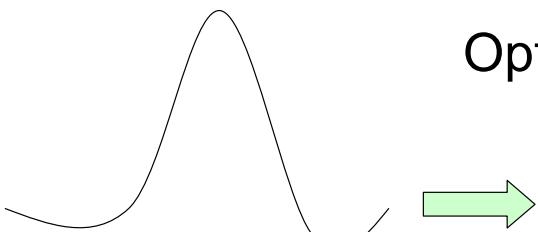
# Desktop Narrow-line MW THz Free-electron Laser



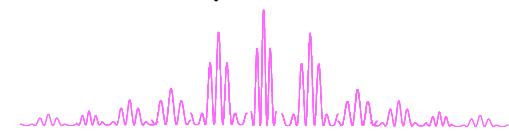
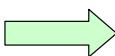
# THz-pulse-train Laser System



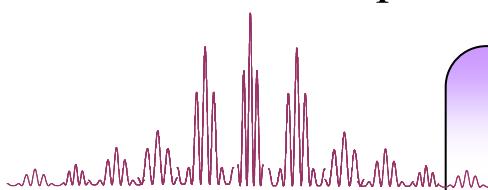
Mode-locked pump at 1064 nm



Optical Parametric Amplifier

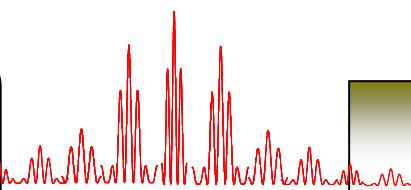


CW, low-power seed laser at 1.56  $\mu\text{m}$   
combined from two diode lasers beating at  
THz frequencies



Ti:sapphire laser  
amplifier + THG

UV laser beat  
wave at 260 nm



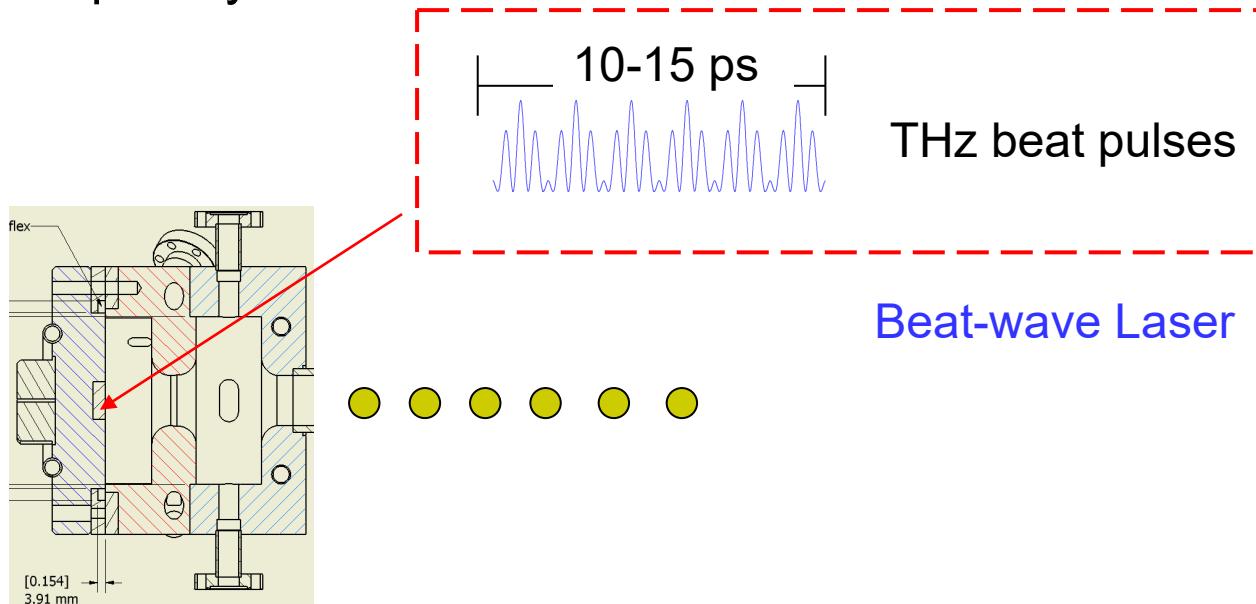
Pulsed laser  
beat wave at  
780 nm

Second harmonic  
generator

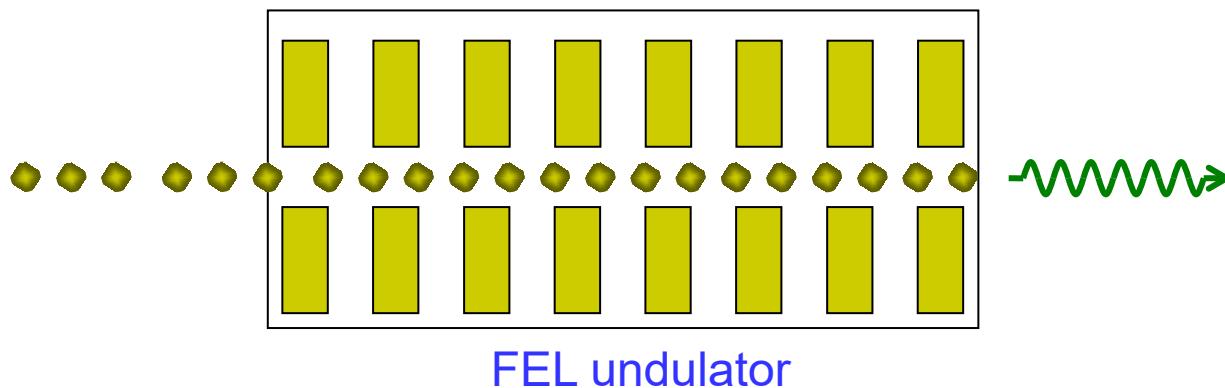
# Simulation Tools



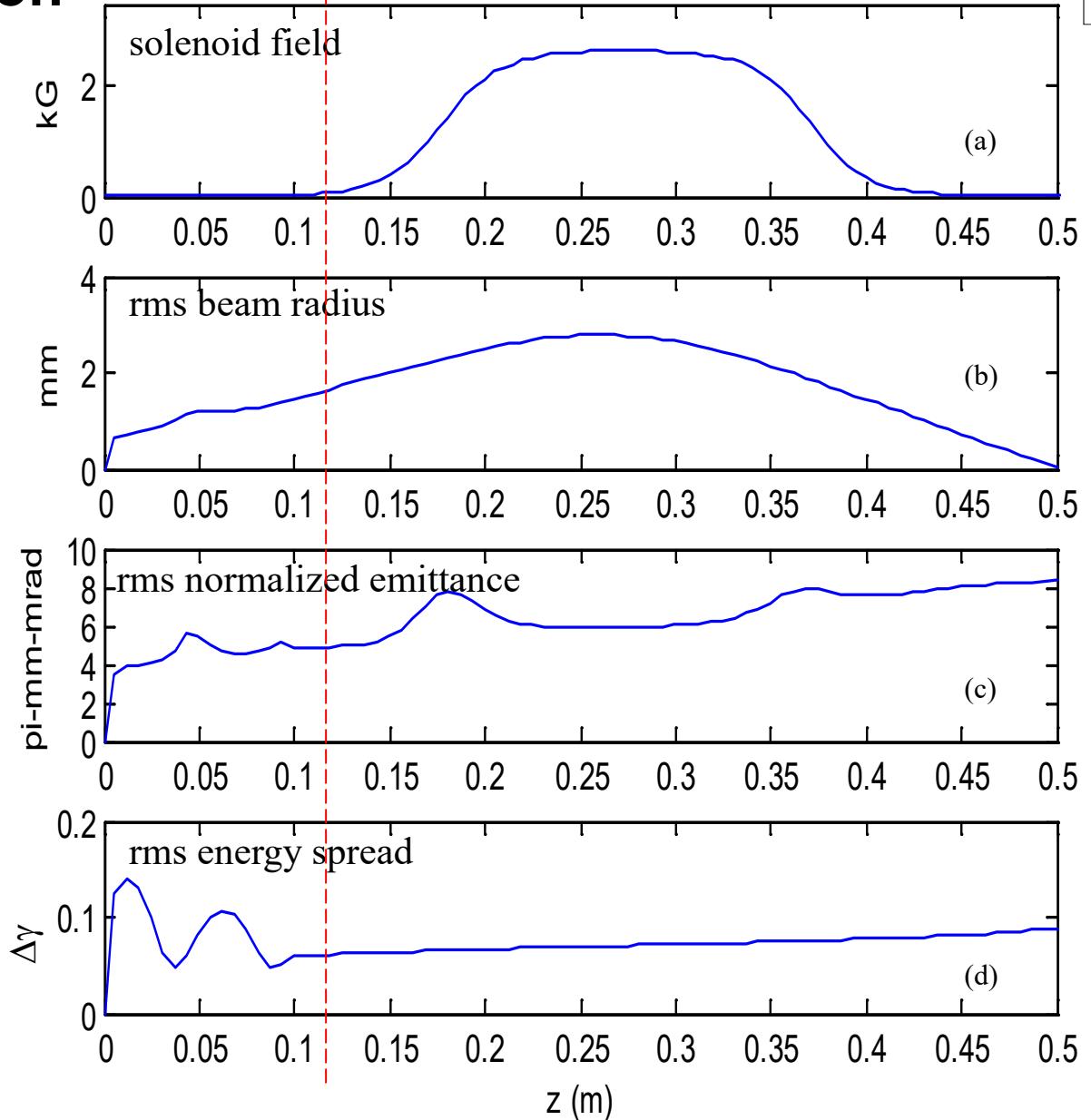
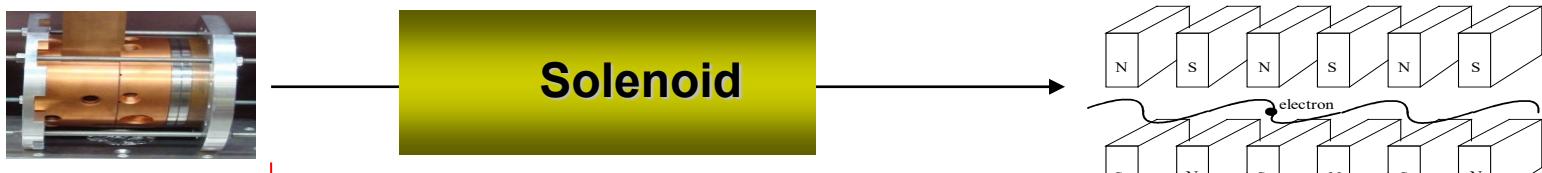
1. Injector Simulation: ASTRA – A space-charge tracking code developed by DESY



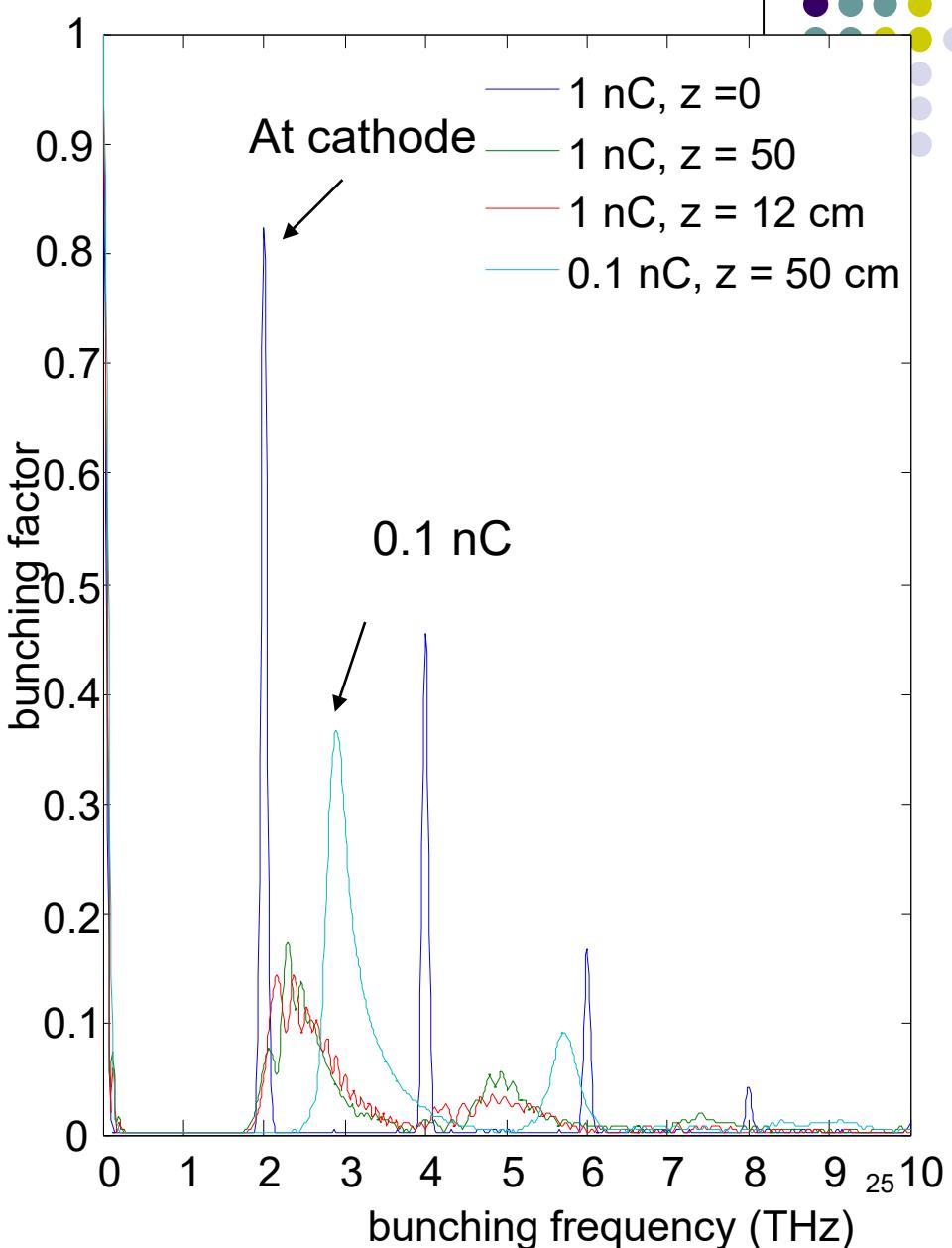
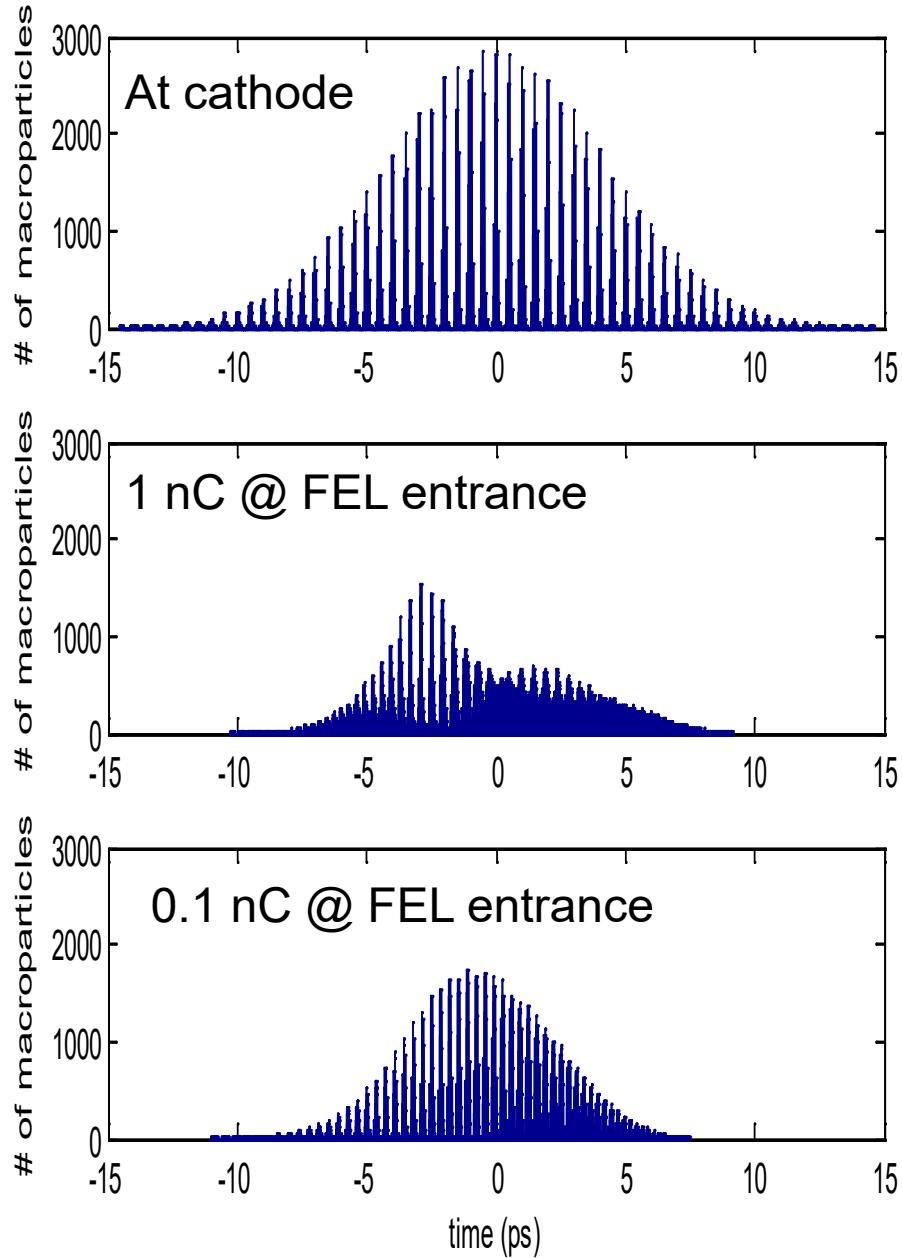
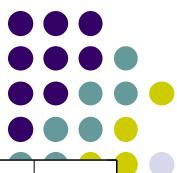
2. FEL Simulation: GENESIS – a SASE FEL code developed by UCLA/DESY



# ASTRA Simulation



# 120 MV/m Acceleration Gradient, comb-f Beat Wave





# Undulator Parameters

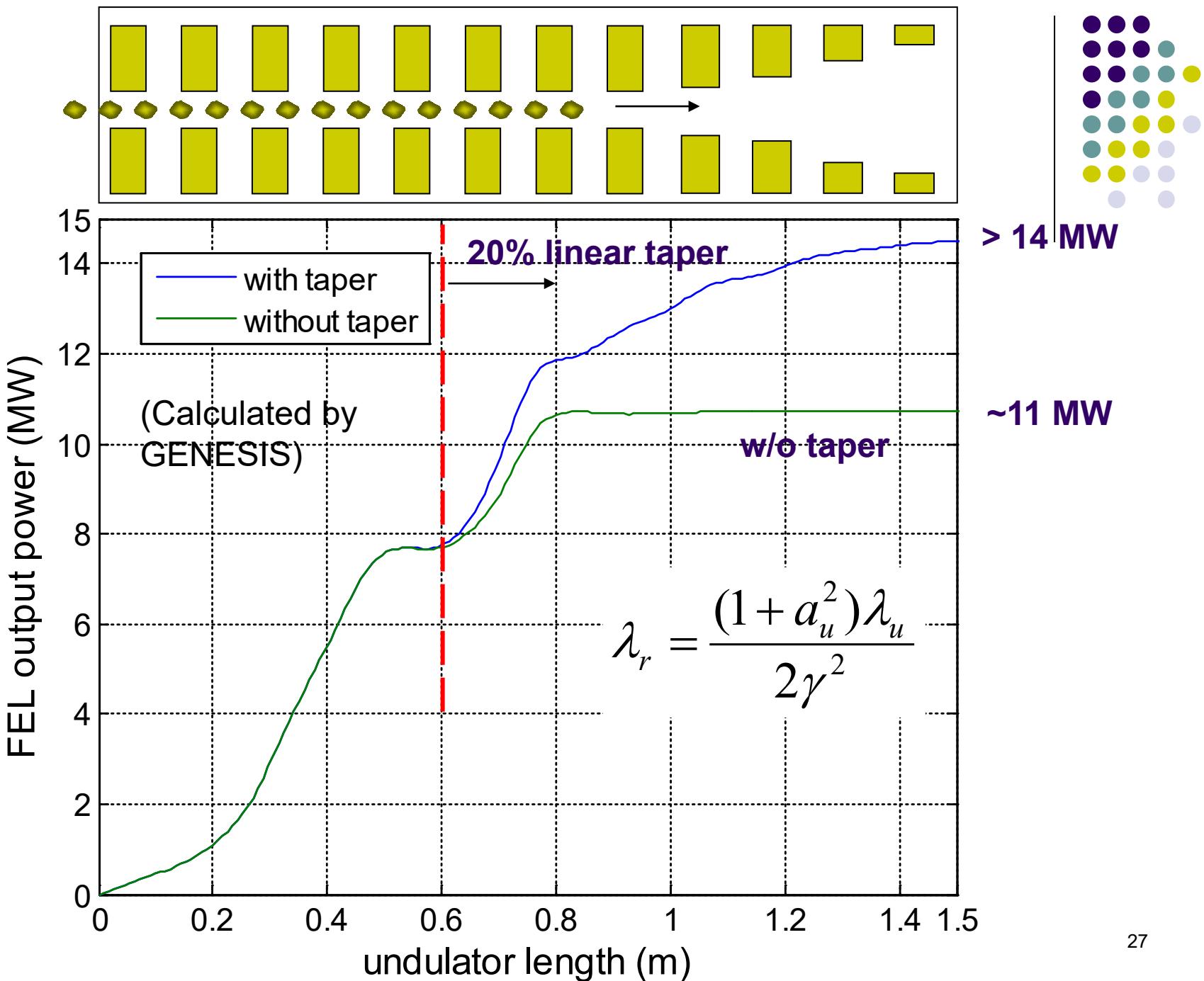
Type: helical

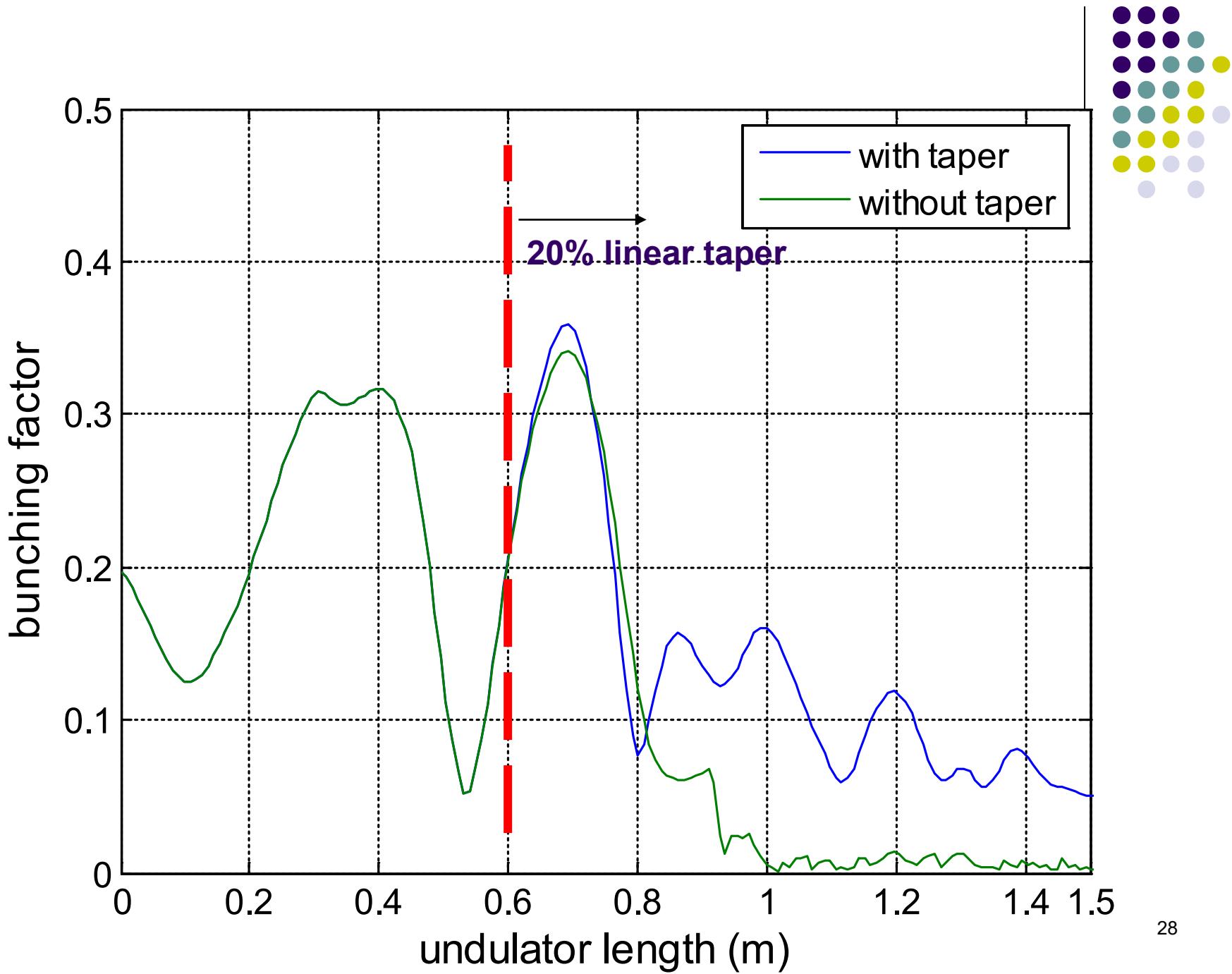
Period: 1.8 cm

Undulator parameter = 0.98

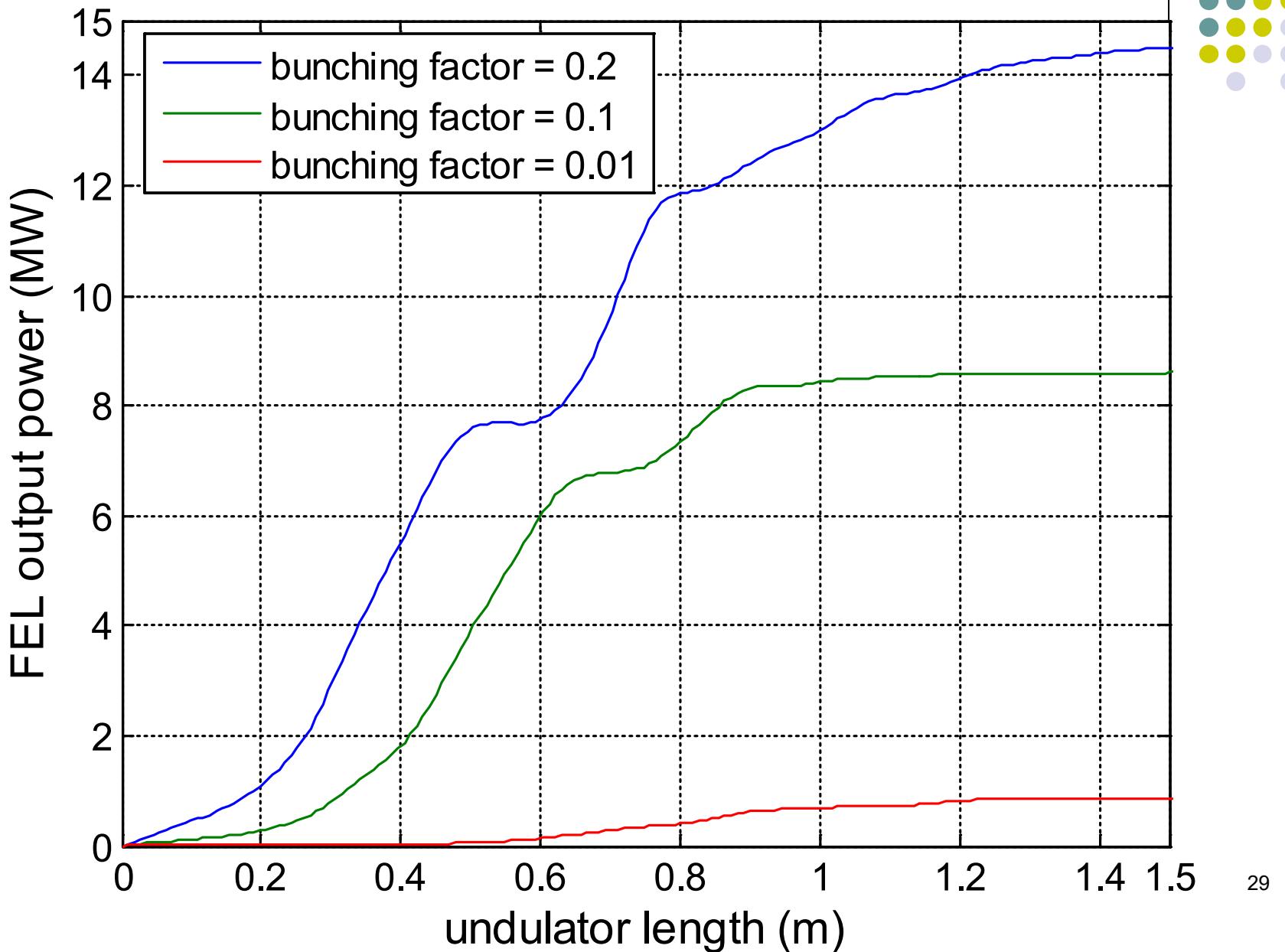
## Beam Parameters (from ASTRA code)

Beam parameters at Cathode						
peak gradient	Charges	rms beam radius (mm)	Macro-bunch length ( $\sigma_M$ )	micro-pulse length ( $\sigma_\mu$ )	micro-pulse rate	Bunching factor @ 2 THz
120 MV/m	1 nC	0.6 (radial distribution)	4.25 ps (10-ps FWHM)	50 fs	2 THz	0.85
Beam Parameters at FEL Entrance						
rms beam energy ( $\gamma$ )	rms energy spread $\Delta\gamma$		rms beam radius (mm)	rms emittance ( $10^{-6} \pi\text{-m}$ )		bunching factor @ 2.4 THz
11.9	$7.4 \times 10^{-2}$ (0.61%)		$6.3 \times 10^{-2}$	6.88		0.21





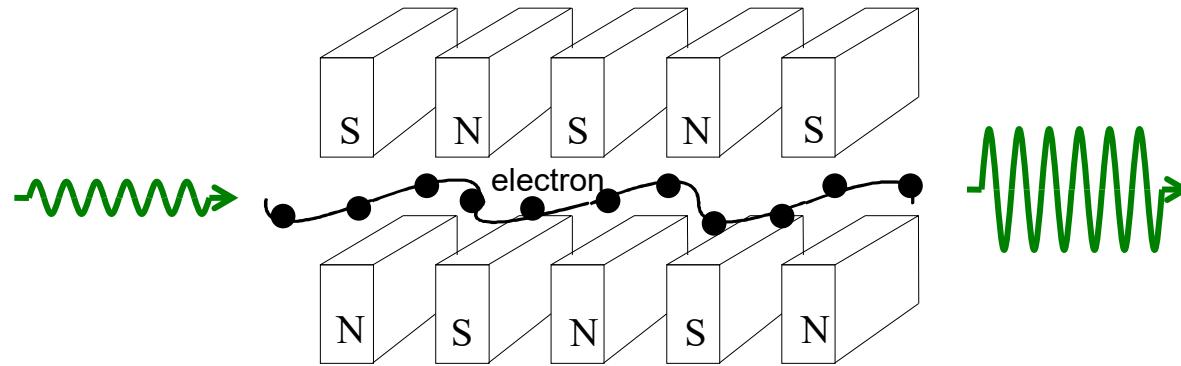
# How critical is the prebunching?



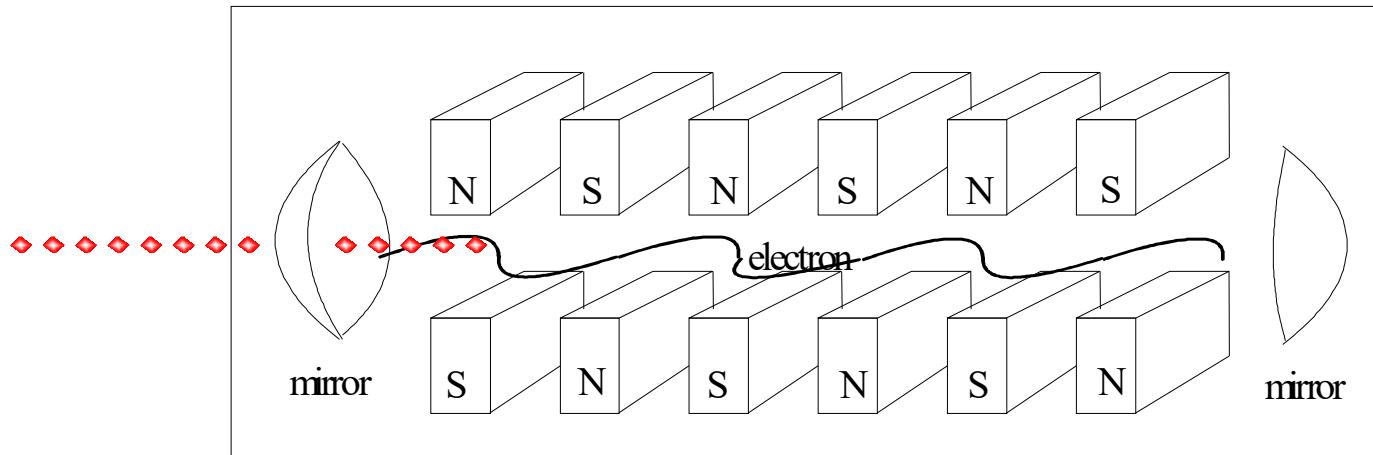
# Variants of Superradiance FEL



## 1. Superradiance FEL amplifier: narrow-line, fully coherent radiation



## 2. Superradiance FEL Oscillator?



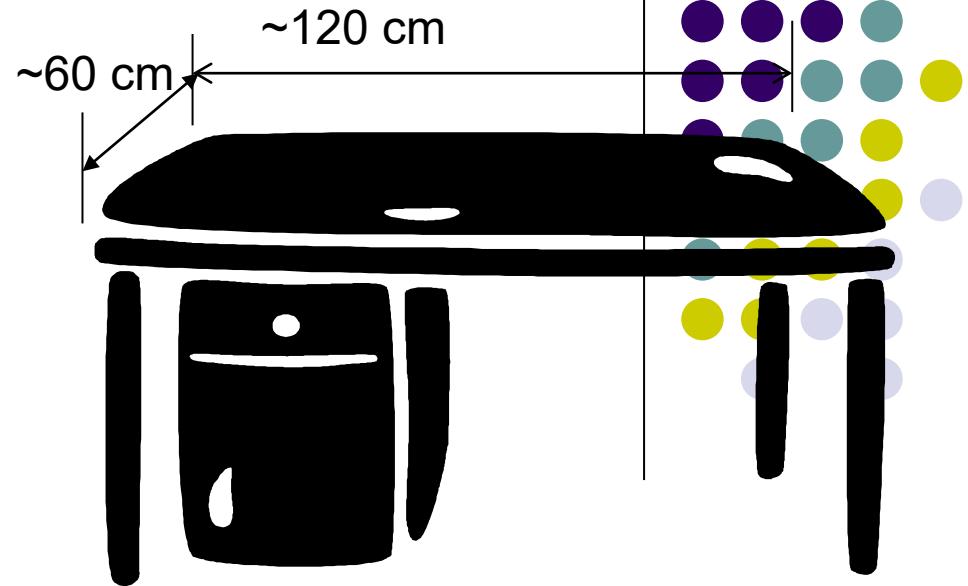
*New physics yet to be discovered*

# “Desktop” MW Superradiant Free-electron Laser at THz Frequencies

High-energy OPtics & Electronics (HOPE) Laboratory  
National Tsinghua University, Hsinchu, Taiwan



**NTHU HOPE Lab**  
(established in Feb. 2008)





## Concept Review

1. A superradiance FEL is driven by well pre-bunched electrons
2. Radiation enters the coherent regime right away with a pre-bunched electron beam.
3. A superradiance FEL skips the slow bunching process in a long undulator, permitting an ultra-compact size.
4. Comb-pulse superradiance narrows down the radiation linewidth.
5. A tapered undulator is recommended for a highly efficiency and ultracompact superradiance FEL.
6. The physics of a superradiance FEL **oscillator** is yet to be investigated.